The Crystal Structure of Silver Chromate

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Single crystals of silver chromate were prepared by slow diffusion of potassium chromate and silver nitrate solutions in a gel-filled U-tube. The thin tabular crystals belong to the orthorhombic space group *Pnma* with a = 10.063(11), b = 7.029(4), c = 5.540(2) Å, and Z = 4. Intensity data were measured with a $\theta - 2\theta$ step-scan technique using an automated diffractometer and Mo K_a radiation. The structure was solved by heavy atom methods and refined by full matrix least-squares techniques with anisotropic thermal parameters used only for the heavy atoms to a final residual of R = 0.064 for the 325 observed reflections used in the analysis. The silver atoms are coordinated to the oxygen atoms of the chromate to generate an interlocking network that propagates throughout the crystal. The two crystallographically independent silver atoms have coordinations best described as tetragonal bipyramidal and distorted tetrahedral, respectively. The shortest Ag–O distance is 2.34 Å. The silver–oxygen interactions perturb the tetrahedron of the chromate group with Cr–O distances of 1.69, 1.67, and 1.63(1) Å and angles varying from 106.1 to 111.6(8)°. The metal atom coordination in silver chromate is compared to that of potassium chromate.

Introduction

During recent studies on alternate methods of growing single crystals, we were successful in growing single crystals of silver chromate. Examination of these crystals by X-ray film techniques revealed the possible orthorhombic space group Pnma. Potassium chromate is also known to crystallize in space group Pnma (1) in a cell of similar dimensions (matching similar unit cell dimensions the equivalent potassium chromate space group would be Pmcn). Although the ionic radii of potassium and silver are similar (1.33 vs 1.26)(2), these compounds exhibit marked differences in properties, such as solubility. Therefore, we felt it would be of interest to proceed with the crystal structure analysis of silver chromate in order to compare the structures of the two compounds.

Experimental

Crystal Data:

Ag₂CrO₄, M = 331.73. Orthorhombic *Pnma*, F(000) = 600 e,

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a = 10.063(11), b = 7.029(4), and c = 5.540(2) Åat $T \approx 24^{\circ}\text{C}$,

Z = 4, V = 391.9 Å³, $D_c = 5.62$ g/cm³, $D_m = 5.62$ g/cm³

Mo
$$K_{\alpha}$$
 ($\lambda = 0.7107$ Å), $\mu = 123$ cm⁻¹.

Thin tabular single crystals of silver chromate were prepared by slow diffusion of potassium chromate and silver nitrate solutions (0.5 M) in a U-tube filled with gel. Precession and Weissenberg photographs exhibited mmm Laue symmetry with the following systematic absences: 0kl when k+l=2n+1 and hk0 when h=2n+1. These extinctions are consistent with either space group *Pn2*, a or *Pnma*. The unit cell parameters and their standard deviations were obtained by a leastsquares fit to 13 independent high-order reflection angles whose centers were determined by left-right, top-bottom beam-splitting techniques on a previously aligned Hilger-Watts four-circle automated diffractometer (Mo K_{α} , $\lambda = 0.71069$ Å). The observed density was taken from the literature (3).

A crystal of approximate dimensions $0.006 \times 0.100 \times 0.118$ mm was mounted on a glass fiber with the *b* axis (0.118 mm) nearly parallel to the ϕ axis and was used for data collection. Data were

taken at room temperature utilizing a fully automated Hilger-Watts four-circle diffractometer equipped with scintillation counter. Mo K_{α} ($\lambda =$ 0.7107 Å) radiation was used with balanced (Zr-Y) filters to obtain intensity and individual background readings. The θ -2 θ step-scan technique with a 4.5° take-off angle was used to record 1056 independent reflections within a 2 θ sphere of 70° (sin $\theta/\lambda = 0.807$ Å⁻¹). A counting rate of 0.2048 sec/step of 0.01° in θ was employed with a variable scan range of 50 steps plus $2/\deg\theta$. The rest of the experimental arrangement has already been discussed in some detail (4). As a general check on electronic and crystal stability, the intensities of three standard reflections were measured periodically during the data collection. Comparison of these values indicated that no decomposition had occurred.

The intensity data were also corrected for Lorentz-polarization effects and for effects due to absorption ($\mu = 123 \text{ cm}^{-1}$). The absorption correction was made using the program of Wehe, Busing, and Levy (5); the maximum and minimum transmission factors were 0.927 and 0.494, respectively. The estimated error in each intensity measurement was calculated by

$$[\sigma(I)]^2 = [C_{\rm T} + C_{\rm B} + (0.05C_{\rm T})^2 + (0.10C_{\rm B})^2 + (0.05C_{\rm R})^2]/A^2,$$

where C_T , C_B , C_R , and A are the total count, background count, net count, and transmission factor, respectively. The estimated standard deviation of each structure factor was calculated from the estimated errors in the intensity by the method of finite differences (6). Based on the measurement of symmetry extinct reflections, it was decided that only those reflections for which $I \ge 3\sigma(I)$ would be considered observed. The results reported below are based on the 325 remaining observed reflections. A very low percentage of observed reflections were obtained at the high 2θ limit.

Solution and Refinement

The positions of all heavy atoms were readily determined from a three-dimensional Patterson function. These positional parameters along with the assigned variable isotropic thermal parameters, were refined two cycles in the centric space group *Pnma* by a full matrix least-squares procedure minimizing the function $\Sigma \omega (|F_o| - |F_c|)^2$. Examination of the resulting electron density map revealed the oxygen positions with the chromate group lying on a mirror plane as required by the centric space group *Pnma*. Isotropic refinement of 18 scale, positional, and

thermal parameters converged to a conventional R factor $(R = (\Sigma ||F_o| - |F_c||/\Sigma |F_o|)$ of 0.082 and to a weighted R-factor

$$R_{\omega} = (\Sigma \omega (|F_o| - |F_c|)^2 / \Sigma \omega |F_o|^2)^{1/2}$$

of 0.090. A difference electron-density map revealed anisotropic motion of the heavy atoms. Additional refinement using anisotropic thermal parameters for the heavy atoms converged to a discrepancy factor of R = 0.064 and a weighted discrepancy factor $R_{\omega} = 0.073$. The largest shift during the last cycle was less than 0.002 times its own standard deviation. A final electron-density difference map showed no peak heights greater than 0.6 $e/Å^3$.

The relativistic Hartree-Fock X-ray scattering factors for neutral atoms of Doyle and Turner (7) were used, with those of silver and chromium including the real and imaginary parts of anomalous dispersion (8). Based on the agreement of the large structure factors no extinction correction was necessary. Anisotropic refinement of the oxygen positions lowered the value of R less than 0.002 with no appreciable shifts in the other parameters. Refinement in the acentric space group produced no significant improvement in either the R factor or standard deviations.

The final positional and thermal parameters and their estimated standard deviations as derived from the inverse matrix of the final least-squares cycle are given in Table I. The final values of the observed

TABLE I

Final Positional and Thermal Parameters for Silver $$\operatorname{Chromate}^{\alpha}$$

Atom	x	ŗ	,	Z	β_{11}^{b} or B
Ag(1)	0.5 ^c	0.5°		0.5°	50(2)
Ag(2)	0.1340(2	2) 0.25°		0.5033(7)	81(2)
Cr	0.3161(3	b) 0.25°		-0.0161(12)	22(2)
O(1)	0.1506(1	9) 0.25 ^c		-0.0677(34)	1.9(4)
O(2)	0.3502(1	9) 0.25 ^c		0.2789(37)	1.8(4)
O(3)	0.3845(1	4) 0.935	3(23)	-0.1417(26)	2.1(2)
Atom	β22	β33	β_{12}	β13	β23
Ag(1)	111(4)	154(5)	Oc	1(7)	Oc
Ag(2)	128(5)	167 (7)	Oc	30(7)	Oc
Cr	81(7)	77(12)	Oc	19(7)	O ^c

^a Numbers in parentheses here and in succeeding tables are estimated standard deviations in the last significant digits.

^b β 's × 10⁴, the β_{ij} are defined by: $T = \exp[-(h^2\beta_{11} + k^2\beta_{22} + l^2\beta_{33} + 2hk\beta_{12} + 2hl\beta_{13} + 2kl\beta_{23}).$

^c Atomic parameter fixed by symmetry.

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TABLE II

Observed and Calculated Structure Factors ($\times 10$) for Silver Chromate

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FIG. 1. Stereogram of the unit cell showing heavy atom coordination polyhedra.

and calculated structure factors ($\times 10$) are listed in Table II. Bond lengths and angles with their standard deviations were calculated by OR FFE. A stereogram of the unit cell is given in Fig. 1.

Description of the Structure

The structure of silver chromate consists of chromate groups linked together by silver atoms to form a three-dimensional network through the crystal. The two crystallographically independent silver atoms have different coordination polyhedra. Positioned on a center of symmetry, the coordination of Ag(1) can be best described as a tetragonal bipyramid with approximate D_{4h} symmetry. The coordination of Ag(2), which lies on the mirror plane, forms a badly distorted tetrahedron with average O-Ag-O angle being 107° but ranging from 84 to 139°. Also on the mirror plane is the chromate group which describes a slightly distorted tetrahedron with average Cr-O distance of 1.66 Å and angles ranging from 106.1 to 111.6(8)°. The coordination of oxygens O(1) and O(2) is roughly tetrahedral with each coordinated to three silver atoms with bond angles of 104 to 125°. The other independent oxygen atom O(3) is coordinated to just two silvers, and the Cr-O(3) distance is 0.04 Å shorter than the remaining Cr-O distances. Bond distances and angles for all the heavy atom coordination polyhedra are listed in Table III (see also Fig. 1). The covalent nature of the silver-oxygen bonds is evidenced by their directional character and comparatively short metal-oxygen distances. Literature values for average Ag-O distances for six coordinate silver (9) and for four coordinate Cr (VI)-O (10) are 2.42 and 1.65 Å, respectively.

Potassium chromate also crystallizes in space group Pnma with Z = 4. Although the unit cell

dimensions are similar, the two structures are not isomorphous; the equivalent potassium chromate space group is *Pmcn*. The metal-atom coordination polyhedra are very different in the two structures. In the ionic potassium salts (K_2CrO_4 and K_2MnO_4) (1) all groups lie on the mirror planes. The potassiumoxygen coordination is complex with little or no directional character and exhibits nine and tenfold coordination. Metal-oxygen distances range from 2.7 to about 3.2 Å compared to the 2.34 to 2.62 Ag-O distances found in silver chromate.

The structural results on these two compounds are in good agreement with their solubilities. Potassium chromate is water soluble and the complex metal coordination found in its structure is typical of ionic

TABLE III

INTERATOMIC DISTANCES AND ANGLES IN SILVER CHROMATE"

	Distance (Å)		Angle (deg)
Ag(1)-O(1)	2.35(1)	O(1)-Ag(1)-O(2)	86.8(4)
Ag(1) - O(2)	2.62(1)	O(1) - Ag(1) - O(3)	87.7(4)
Ag(1)–O(3)	2.34(1)	O(2)-Ag(1)-O(3)	88.9(4)
Ag(2)–O(1)	2.38(2)	O(1)-Ag(2)-O(2)	115.7(6)
Ag(2) - O(2)	2.51(2)	O(1) - Ag(2) - O(3)	110.2(3)
Ag(2) - O(3)	2.36(2)	O(2) - Ag(2) - O(3)	84.3(4)
		O(3)-Ag(2)-O(3)	139.1(7)
Cr-O(1)	1.69(1)	O(1)-Cr-O(2)	111.6(9)
CrO(2)	1.67(1)	O(1)-Cr-O(3)	110.1(6)
Cr-O(3)	1.63(1)	O(2)-Cr-O(3)	109.3(6)
		O(3)-Cr-O(3)	106.1(10)

^a Symbols used refer to atom type only (Fig. 1). Symmetry operations are not given. Ag(1) lies on a center of symmetry so all other O-Ag(1)-O angles are related to those given.

compounds. Silver chromate is quite insoluble in polar solvents such as water and the more directional metal coordination of its structure is more representative of covalent compounds. Also, the electronegativity of silver (1.42) (11) is close to that of chromium (1.56) which is dissimilar to that of potassium (0.91). Thus the ionic compound potassium chromate crystallizes with complex potassiumoxygen coordinations providing many near neighbors to minimize the electrostatic repulsions of the ions, while in the more covalent silver chromate there is substantial directional silver-oxygen coordination to provide maximum overlap for some formation of covalent bonds. The differences in the two structures is, therefore, in accord with the different nature of the binding forces for the two solids.

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